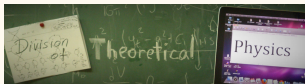


# SIGMA-MODELS & NON-GEOMETRIC FLUXES IN STRING AND M THEORIES

Athanasios Chatzistavrakidis



Rudjer Bošković Institute, Zagreb



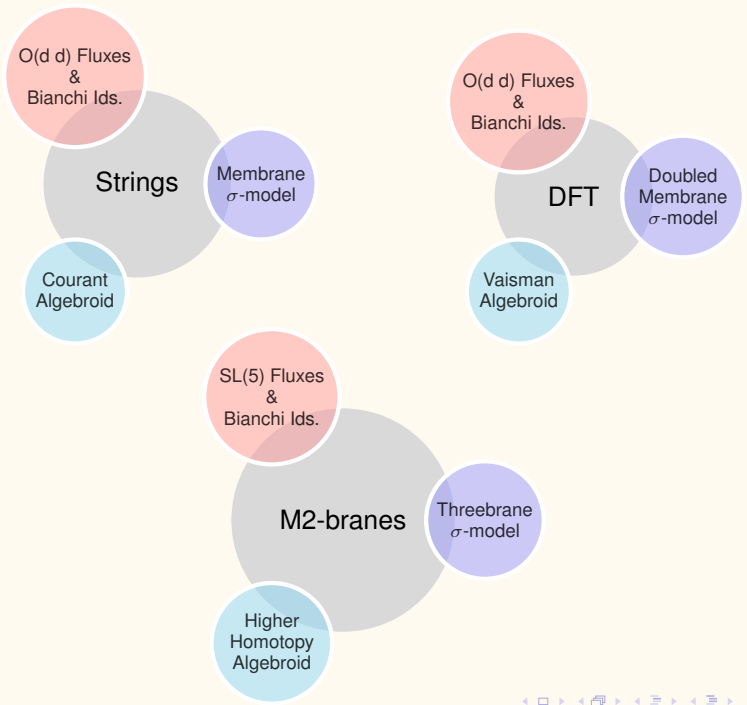
Based on:

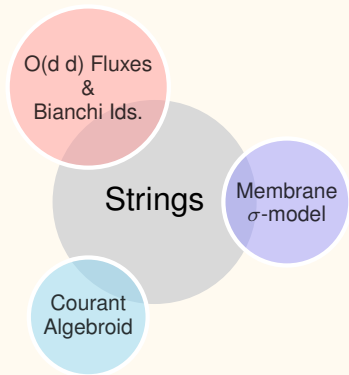
1901.07775 with L. Jonke - D. Lüst - R. J. Szabo

1802.07003 with L. Jonke - F. S. Khoo - R. J. Szabo

1311.4878 & 1505.05457 with L. Jonke - O. Lechtenfeld

Quantum Spacetime '19





## Three equations, three viewpoints

$$\eta^{IJ} \rho^i{}_I \rho^j{}_J = 0, \quad i = 1, \dots, d$$

$$\rho^i{}_I \partial_i \rho^j{}_J - \rho^i{}_J \partial_i \rho^j{}_I - \eta^{KL} \rho^j{}_K T_{LJ} = 0, \quad I = 1, \dots, 2d$$

$$4 \rho^i{}_{[L} \partial_i T_{JK]} + 3 \eta^{MN} T_{M[L} T_{KL]N} = 0.$$

$\eta$  is the  $O(d, d)$ -invariant metric:  $(\eta_{IJ}) = \begin{pmatrix} 0 & \mathbf{1} \\ \mathbf{1} & 0 \end{pmatrix}$

## Viewpoint 1: Flux compactifications of strings

$\rho$ : Potential for geometric and non-geometric fluxes  $\rightsquigarrow \rho^i{}_l = (e^i{}_a, e^i{}_b \beta^{ba})$

$T$ : The (NSNS) fluxes  $\rightsquigarrow T_{IJK} = (H_{abc}, f_{ab}{}^c, Q_a{}^{bc}, R^{abc})$ , typically related by T-duality

$$H_{abc} \xleftrightarrow{T_c} f_{ab}{}^c \xleftrightarrow{T_b} Q_a{}^{bc} \xleftrightarrow{T_a} R^{abc} .$$

Decomposing the second equation, one obtains the local expressions for all fluxes:

$$H_{abc} = 3 \nabla_{[a} B_{bc]},$$

$$F_{ab}{}^c = 2 e^c{}_j e_{[a}{}^i \partial_i e_{b]}{}^j + \beta^{cd} H_{abd} ,$$

$$Q_a{}^{bc} = \partial_a \beta^{bc} + \beta^{bd} f_{ad}{}^c - \beta^{cd} f_{ad}{}^b + \beta^{bd} \beta^{ce} H_{ade} ,$$

$$R^{abc} = 3 \beta^{[ad} \nabla_d \beta^{bc]} + \beta^{ad} \beta^{be} \beta^{cf} H_{def} ,$$

Decomposing the third equation, one obtains their Bianchi identities.

## Viewpoint 2: Gauge structure of membrane sigma models

Courant sigma-model for maps  $X = (X^i) : \Sigma_3 \rightarrow M_d$  with Wess-Zumino term  $T$

Hofman, Park '02; Ikeda '02

$$S_0[X, A, F] = \int_{\Sigma_3} (F_i \wedge dX^i + \frac{1}{2} \eta_{IJ} A^I \wedge dA^J - \rho^i{}_I(X) A^I \wedge F_i + \frac{1}{6} T_{IJK}(X) A^I \wedge A^J \wedge A^K) .$$

The 1-forms  $A$  take values in the (pull-back) bundle  $E = TM \oplus T^*M$ .

The 3 equations encode the gauge symmetries of the theory & their on-shell closure

$$\begin{aligned} \delta_\epsilon X^i &= \rho^i{}_J \epsilon^J , \\ \delta_\epsilon A^I &= d\epsilon^I + \eta^{IN} T_{NJK} A^J \epsilon^K - \eta^{IJ} \rho^i{}_J t_i , \\ \delta_\epsilon F_m &= -dt_m - \partial_m \rho^j{}_J A^J t_j - \epsilon^J \partial_m \rho^i{}_J F_i + \frac{1}{2} \epsilon^J \partial_m T_{ILJ} A^I \wedge A^L . \end{aligned}$$

Equivalently, the classical master equation of the corresponding BV action is satisfied.

On-shell (w/ suitable boundary kinetic term)  $\rightsquigarrow$  strings in (non-)geometric backgrounds.

Mylonas, Schupp, Szabo '12; ACh, Jonke, Lechtenfeld '15

## Viewpoint 3: Axioms of a Courant algebroid

Axiomatic organisation of the properties of the (twisted) Courant bracket on  $TM \oplus T^*M$

Courant '90; Liu, Weinstein, Xu '95; Ševera

$$[X + \xi, Y + \eta] = [X, Y] + \mathcal{L}_X \eta - \mathcal{L}_Y \xi - \frac{1}{2} d(\iota_X \eta - \iota_Y \xi) + H(X, Y).$$

$(E \xrightarrow{\pi} M, [\cdot, \cdot], \langle \cdot, \cdot \rangle, \rho : E \rightarrow TM)$ , such that for  $A, B, C \in \Gamma(E)$  and  $f, g \in C^\infty(M)$ :

- 1  $[[A, B], C] + \text{c.p.} = \mathcal{DN}(A, B, C)$ , where  $\mathcal{N}(A, B, C) = \frac{1}{3} \langle [A, B], C \rangle + \text{c.p.}$ ,
- 2  $[A, fB] = f[A, B] + (\rho(A)f)B - \langle A, B \rangle \mathcal{D}f$ ,
- 3  $\rho(C) \langle A, B \rangle = \langle [C, A] + \mathcal{D} \langle C, A \rangle, B \rangle + \langle [C, B] + \mathcal{D} \langle C, B \rangle, A \rangle$ ,
- 4  $\rho[A, B] = [\rho(A), \rho(B)]$ ,
- 5  $\rho \circ \mathcal{D} = 0 \Leftrightarrow \langle \mathcal{D}f, \mathcal{D}g \rangle = 0$ .

In local coordinates, they become the three equations...

O(d d) Fluxes  
&  
Bianchi Ids.

Double Field Theory

Doubled  
Membrane  
 $\sigma$ -model

Vaisman  
Algebroid



## DFT: Three equations, three viewpoints

$$\eta^{IJ} \rho^K{}_I \rho^L{}_J = \eta^{KL}$$

$$2\rho^L{}_{[I} \partial_L \rho^K{}_{J]} - \eta^{LM} \rho^K{}_L \hat{T}_{MJ} = \rho_{L[I} \partial^K \rho^L{}_{J]}$$

$$4\rho^M{}_{[L} \partial_M \hat{T}_{JK]} + 3\eta^{MN} \hat{T}_{M[J} \hat{T}_{KL]N} = \mathcal{Z}_{JKL} .$$

## Viewpoint 1: Duality-symmetric formulation of fluxes

$\rho$  and  $\hat{T}$  become a generalized vielbein and the T-duality-symmetric DFT fluxes.

Decomposing the second equation, one obtains the augmented local expressions:

Aldazabal, Baron, Marques, Nunez '11; Geissbühler '11

$$H_{ijk} = 3 \partial_{[i} B_{jk]} + 3 B_{[i|} \tilde{\partial}^j B_{jk]} ,$$

$$f_{ij}{}^k = \tilde{\partial}^k B_{ij} + \beta^{kl} H_{lij} ,$$

$$Q_k{}^{ij} = \partial_k \beta^{ij} + B_{kl} \tilde{\partial}^l \beta^{ij} + 2 \beta^{[i} \tilde{\partial}^{j]} B_{lk} + \beta^{il} \beta^{jm} H_{lmk} ,$$

$$R^{ijk} = 3 \tilde{\partial}^{[i} \beta^{jk]} + 3 \beta^{[i|} \partial_l \beta^{jk]} \\ + 3 B_{lm} \beta^{[i|} \tilde{\partial}^m \beta^{jk]} + 3 \beta^{[i|} \beta^{jm} \tilde{\partial}^k] B_{lm} + \beta^{il} \beta^{jm} \beta^{kn} H_{lmn} ,$$

Decomposing the third equation, one obtains their corresponding Bianchi identities.

## Viewpoint 2: Doubled membrane sigma-model

A double worth of scalar fields  $(\mathbb{X}^I) = (X^i, \tilde{X}_i)$ : ACh, Jonke, Lechtenfeld '15; ACh, Jonke, Khoo, Szabo '18

$$S[\mathbb{X}, \mathbf{A}, F] = \int \left( F_I \wedge d\mathbb{X}^I + \eta_{IJ} A^I \wedge dA^J - (\rho_+)^I{}_J A^J \wedge F_I + \frac{1}{3} \hat{T}_{IJK} A^I \wedge A^J \wedge A^K \right) .$$

Gauge symmetries and their closure (or the classical master equation), obstructed.

Obstructions vanish when a strong constraint is satisfied:  $\eta^{IJ} \partial_I \otimes \partial_J = 0$  . cf. Hull, Zwiebach '09

On-shell, describes non-commutative/non-associative string flux backgrounds.

## Viewpoint 3: (Pre-)DFT (Vaisman) Algebroid

Axiomatic organisation of the properties of the C-bracket: Siegel '93; Hull, Zwiebach '09

$$[[A, B]]^J = A^K \partial_K B^J - \frac{1}{2} A^K \partial^J B_K - \{A \leftrightarrow B\} .$$

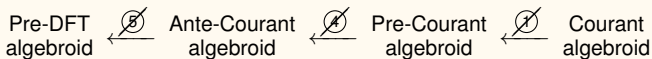
$(L_+ \xrightarrow{\pi} \mathcal{M}, [[\cdot, \cdot]], \langle \cdot, \cdot \rangle_{L_+}, \rho_+ : L_+ \rightarrow T\mathcal{M})$  satisfying

- ②  $[[A, fB]] = f[[A, B]] + (\rho_+(A)f)B - \langle A, B \rangle_{L_+} \mathcal{D}_+ f ,$
- ③  $\langle [[C, A]] + \mathcal{D}_+ \langle C, A \rangle_{L_+}, B \rangle_{L_+} + \langle [[C, B]] + \mathcal{D}_+ \langle C, B \rangle_{L_+}, A \rangle_{L_+} = \rho_+(C) \langle A, B \rangle_{L_+} ,$
- ⑤  $\langle \mathcal{D}_+ f, \mathcal{D}_+ g \rangle_{L_+} = \frac{1}{4} \langle df, dg \rangle_{L_+} .$

Notably, the modified Jacobi, homomorphism and kernel properties are obstructed.

In general, by relaxing properties one obtains a host of intermediate structures

cf. Vaisman '04; Hansen, Strobl '09; Vaisman '12; Bruce, Grabowski '16



SL(5) Fluxes  
&  
Bianchi Ids.

M2-branes

Threebrane  
 $\sigma$ -model

Higher  
Homotopy  
Algebroid

## Three + Two equations, three viewpoints

$$\rho^i{}_I S^{IJ} = 0 ,$$

$$\rho^i{}_I \partial_i S^{JK} + S^{LJ} T^K{}_{IL} + S^{LK} T^J{}_{IL} = 0 ,$$

$$\rho^i{}_I \partial_i \rho^j{}_J - \rho^i{}_J \partial_i \rho^j{}_I - \rho^i{}_K T^K{}_{IJ} = 0 ,$$

$$3\rho^i{}_{[I} \partial_i T^J{}_{KL]} + S^{JM} G_{KLIM} - 3T^J{}_{M[K} T^M{}_{L]} = 0 ,$$

$$\rho^i{}_{[I} \partial_i G_{JKLM]} + T^N{}_{[IJ} G_{KLM]N} = 0 ,$$

N.B. Index  $I = 1, \dots, \dim E$ , where  $E$  is a vector bundle.

## Viewpoint 1: Non-geometric M theory fluxes

When 11D supergravity is compactified on  $T^4 \rightsquigarrow$  SL(5) (U-duality) symmetry.

Cremmer, Julia '78, '81

Fluxes are irreps of SL(5):  $G_{abcd}$ ,  $F_{ab}{}^c$ ,  $Q_a{}^{bcd}$ ,  $R^{abcde}$ .

Blair, Malek '14

Alternatively, twists of the (higher) Courant bracket on  $TM \oplus \wedge^2 T^*M$  for  $M$  being 4D.

Their general local coordinate expressions are found to be

$$G_{abcd} = 4 \nabla_{[a} C_{bcd]} ,$$

$$F_{ab}{}^c = f_{ab}{}^c - \frac{1}{2} G_{abde} \Omega^{dec} ,$$

$$Q_a{}^{bcd} = \frac{1}{2} (\partial_a \Omega^{bcd} + 3 \Omega^{e[bc} f_{ae}{}^{d]} - \frac{1}{2} \Omega^{def} \delta_a^{[b} f_{ef}{}^{c]} - \frac{1}{2} \Omega^{e[bc} G_{aefg} \Omega^{d]fg}) ,$$

$$R^{ab,cd,e} = 2 \widehat{\nabla}^{a[b} \Omega^{cde]} - 2 \widehat{\nabla}^{b[a} \Omega^{cde]} - 2 \widehat{\nabla}^{c[d} \Omega^{abe]} + 2 \widehat{\nabla}^{d[c} \Omega^{abe]} ,$$

where  $\widehat{\nabla}^{ab} = \frac{1}{4} \Omega^{abc} \nabla_c$  and  $f_{ab}{}^c = 2 e^c{}_j e_{[a}{}^i \partial_i e_{b]}{}^j =: 2 \Gamma_{[ab]}{}^c$  the purely geometric flux.

These and their BIs can actually be directly obtained from the three equations ( $S = 0 \dots$ )

## Viewpoints 2 & 3: Threebrane sigma-model & higher homotopy algebroids

The starting point is is a topological threebrane sigma-model with action functional

Ikeda, Uchino '10; Kökenyesi, Sinkovics, Szabo '18

$$S[X, \alpha, A, F] = \int_{\Sigma_4} (F_i \wedge dX^i - \alpha_I \wedge dA^I + \rho^i{}_I(X) F_i \wedge A^I + \frac{1}{2} S^{IJ}(X) \alpha_I \wedge \alpha_J + \frac{1}{2} T^I{}_{JK}(X) \alpha_I \wedge A^J \wedge A^K + \frac{1}{4!} G_{IJKL}(X) A^I \wedge A^J \wedge A^K \wedge A^L) .$$

Here,  $A \in \Omega^1(\Sigma_4, X^*E)$  &  $\alpha \in \Omega^2(\Sigma_4, X^*E^*)$ . Gauge structure  $\rightsquigarrow$  the five equations.

(Moreover, local form of the axioms for a structure called Lie algebroid up to homotopy.)

Ikeda, Uchino '10; Grützmann '10

Specialize:  $E = TM \oplus \wedge^2 T^*M$  (N.B. Not a standard choice!), then  $T = \text{SL}(5)$  fluxes.



# Hierarchy of structures

<u>dim <math>\Sigma</math></u>	<u><math>\sigma</math>-model</u>	<u>Scalars</u>	<u>1-forms</u>	<u>2-forms</u>	<u>3-forms</u>
2	Poisson	$X^i$	$F_i \in \Gamma(X^* T^* M)$	—	—
3	Courant	$X^i$	$A^I \in \Gamma(X^* E)$	$F_i \in \Gamma(X^* T^* M)$	—
4	Threebrane	$X^i$	$A^I \in \Gamma(X^* E)$	$\alpha_I \in \Gamma(X^* E^*)$	$F_i \in \Gamma(X^* T^* M)$

<u>dim <math>\Sigma</math></u>	<u><math>\sigma</math>-model</u>	<u>Algebroid</u>	<u>Quantization</u>
2	Poisson	Lie	Open string non-commutativity Kontsevich '97; Cattaneo, Felder '99
3	Courant	Courant	Closed string non-associativity Mylonas, Schupp, Szabo '12
4	Threebrane	Homotopy	...