

On generic superconformal field theories on $K3$

QUANTUM STRUCTURE OF SPACETIME

Bratislava, Slovakia, February 11–15, 2019

Katrin Wendland

Albert-Ludwigs-Universität Freiburg

Plan:

- 1 Superconformal field theories on $K3$
- 2 Refining the conformal field theoretic elliptic genus
- 3 Elliptic genera in geometry
- 4 Results

[W17] *Hodge-elliptic genera and how they govern $K3$ theories*, accepted for publication by Comm. Math. Phys.; arXiv:1705.09904 [hep-th]

[Taormina/W15] *A twist in the M_{24} moonshine story*, Confluentes Mathematici 7, 1 (2015), 83–113; arXiv:1303.3221 [hep-th]

[Taormina/W13] *Symmetry-surfing the moduli space of Kummer $K3$ s*, Proceedings of Symposia in Pure Mathematics 90 (2015), 129–153; arXiv:1303.2931 [hep-th]

[Taormina/W11] *The overarching finite symmetry group of Kummer surfaces in the Mathieu group M_{24}* , JHEP **1308**:152 (2013); arXiv:1107.3834 [hep-th]

1. Superconformal field theories on K3 - what they are

M : a compact Calabi-Yau D -fold

parameters of non-linear sigma-models on M :

complex structure, Kähler-Einstein metric, B-field

Result:

- [Casher/Englert/Nicolai/Taormina85, Narain86]

The moduli space \mathcal{M}_{T^D} of non-linear sigma models on a complex D -torus is

$$\mathcal{M}_{T^D} = O(D, D; \mathbb{Z}) \backslash O(D, D; \mathbb{R}) / O(D) \times O(D).$$

- [Seiberg88, Cecotti90; Aspinwall/Morrison94; Nahm/W00]

The moduli space $\overline{\mathcal{M}}_{K3}$ of non-linear sigma models on K3 is

$$\overline{\mathcal{M}}_{K3} = O^+(4, 20; \mathbb{Z}) \backslash O(4, 20; \mathbb{R}) / O(4) \times O(20).$$

1. Superconformal field theories - what we assume

Notations and assumptions:

- \mathbb{H} : Neveu-Schwarz sector of a unitary 2-dimensional $N = (2, 2)$ superconformal field theory, such that
- the spectral flow isomorphically maps \mathbb{H} to the Ramond sector (space-time supersymmetry)
 - the central charges are $c = \bar{c} = 3D$ with $D \in \mathbb{N}$
 - all eigenvalues of J_0 , \bar{J}_0 and $H - \bar{H}$ are integral, where

$$2H := 2L_0 - J_0, \quad 2\bar{H} := 2\bar{L}_0 - \bar{J}_0$$

With $\tau, z \in \mathbb{C}$, $\text{Im}(\tau) > 0$, $q := \exp(2\pi i\tau)$, $y := \exp(2\pi iz)$,

WITTEN GENUS:

$$\chi(\mathbb{H}) := \text{tr}_{\mathbb{H}} \left((-1)^{J_0 - \bar{J}_0} q^{H - \bar{H}} \right)$$

CONFORMAL FIELD THEORETIC ELLIPTIC GENUS:

$$\mathcal{E}_{\text{CFT}}(\mathbb{H}; \tau, z) := \text{tr}_{\mathbb{H}} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} q^{H - \bar{H}} \right)$$

1. Superconformal field theories - what we assume

Notations and assumptions:

- \mathbb{H} : **Neveu-Schwarz sector** of a unitary 2-dimensional $N = (2, 2)$ superconformal field theory, such that
- the **spectral flow isomorphically** maps \mathbb{H} to the **Ramond sector** (space-time supersymmetry)
 - the **central charges** are $c = \bar{c} = 3D$ with $D \in \mathbb{N}$
 - all eigenvalues of J_0 , \bar{J}_0 and $H - \bar{H}$ are **integral**, where

$$2H := 2L_0 - J_0, \quad 2\bar{H} := 2\bar{L}_0 - \bar{J}_0$$

With $\tau, z \in \mathbb{C}$, $\text{Im}(\tau) > 0$, $q := \exp(2\pi i\tau)$, $y := \exp(2\pi iz)$,

WITTEN GENUS:

$$\chi(\mathbb{H}) := \text{tr}_{\mathbb{H}} \left((-1)^{J_0 - \bar{J}_0} q^H \bar{q}^{\bar{H}} \right)$$

CONFORMAL FIELD THEORETIC ELLIPTIC GENUS:

$$\mathcal{E}_{\text{CFT}}(\mathbb{H}; \tau, z) := \text{tr}_{\ker(\bar{H})} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} q^H \bar{q}^{\bar{H}} \right)$$

1. Superconformal field theories - what we assume

Notations and assumptions:

- \mathbb{H} : **Neveu-Schwarz sector** of a unitary 2-dimensional $N = (2, 2)$ **superconformal field theory**, such that
- the **spectral flow isomorphically** maps \mathbb{H} to the **Ramond sector** (space-time supersymmetry)
 - the **central charges** are $c = \bar{c} = 3D$ with $D \in \mathbb{N}$
 - all eigenvalues of J_0 , \bar{J}_0 and $H - \bar{H}$ are **integral**, where

$$2H := 2L_0 - J_0, \quad 2\bar{H} := 2\bar{L}_0 - \bar{J}_0$$

With $\tau, z \in \mathbb{C}$, $\text{Im}(\tau) > 0$, $q := \exp(2\pi i\tau)$, $y := \exp(2\pi iz)$,

WITTEN GENUS:

$$\chi(\mathbb{H}) := \text{tr}_{\mathbb{H}} \left((-1)^{J_0 - \bar{J}_0} q^{H - \bar{H}} \right) = \mathcal{E}_{\text{CFT}}(\mathbb{H}; \tau, z = 0)$$

CONFORMAL FIELD THEORETIC ELLIPTIC GENUS:

$$\mathcal{E}_{\text{CFT}}(\mathbb{H}; \tau, z) := \text{tr}_{\ker(\bar{H})} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} q^{H - \bar{H}} \right)$$

1. Superconformal field theories - what we assume

Notations and assumptions:

- \mathbb{H} : **Neveu-Schwarz sector** of a unitary 2-dimensional $N = (2, 2)$ superconformal field theory, such that
- the **spectral flow isomorphically** maps \mathbb{H} to the **Ramond sector** (space-time supersymmetry)
 - the **central charges** are $c = \bar{c} = 3D$ with $D \in \mathbb{N}$
 - all eigenvalues of J_0 , \bar{J}_0 and $H - \bar{H}$ are **integral**, where

$$2H := 2L_0 - J_0, \quad 2\bar{H} := 2\bar{L}_0 - \bar{J}_0$$

With $\tau, z \in \mathbb{C}$, $\text{Im}(\tau) > 0$, $q := \exp(2\pi i\tau)$, $y := \exp(2\pi iz)$,

WITTEN GENUS:

$$\chi(\mathbb{H}) := \text{tr}_{\mathbb{H}} \left((-1)^{J_0 - \bar{J}_0} q^{H - \bar{H}} \right) = \mathcal{E}_{\text{CFT}}(\mathbb{H}; \tau, z = 0) \in \mathbb{Z}$$

CONFORMAL FIELD THEORETIC ELLIPTIC GENUS:

$$\mathcal{E}_{\text{CFT}}(\mathbb{H}; \tau, z) := \text{tr}_{\ker(\bar{H})} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} q^{H - \bar{H}} \right)$$

1. Defining K3 theories

Definition

A **K3 THEORY** is a superconformal field theory as above at $c = 6$ with **Witten index** $\chi(\mathbb{H}) = 24$.

Results:

- **[Eguchi/Ooguri/Taormina/Yang89]**

All **K3 theories** possess $N = (4, 4)$ supersymmetry and have **conformal field theoretic elliptic genus**

$$\mathcal{E}_{\text{CFT}}(\mathbb{H}; \tau, z) = 8 \left(\frac{\vartheta_2(\tau, z)}{\vartheta_2(\tau, 0)} \right)^2 + 8 \left(\frac{\vartheta_3(\tau, z)}{\vartheta_3(\tau, 0)} \right)^2 + 8 \left(\frac{\vartheta_4(\tau, z)}{\vartheta_4(\tau, 0)} \right)^2.$$

- **[Cecotti90; Aspinwall/Morrison94; Nahm/W00]**

There is an **80-dimensional moduli space** of **K3 theories**, one connected component being \mathcal{M}_{K3} with

$$\overline{\mathcal{M}}_{\text{K3}} = \text{O}^+(4, 20; \mathbb{Z}) \backslash \text{O}(4, 20; \mathbb{R}) / \text{O}(4) \times \text{O}(20).$$

2. The conformal field theoretic Hodge elliptic genera

ELLIPTIC GENUS

$$\mathcal{E}_{\text{CFT}}(\mathbb{H}; \tau, z) = \text{tr}_{\ker(\bar{H})} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} q^H \right)$$

2. The conformal field theoretic Hodge elliptic genera

HODGE ELLIPTIC GENERA

With $\nu \in \mathbb{C}$, $u := \exp(2\pi i\nu)$,

- [Kachru/Tripathy16]

$$\mathcal{E}_{\text{CFT}}^{\text{HEG}}(\mathbb{H}; \tau, z, \nu) := \text{tr}_{\ker(\bar{H})} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} u^{\bar{J}_0 - c/6} q^H \right).$$

2. The conformal field theoretic Hodge elliptic genera

HODGE ELLIPTIC GENERA

With $\nu \in \mathbb{C}$, $u := \exp(2\pi i\nu)$,

- [Kachru/Tripathy16]

$$\mathcal{E}_{\text{CFT}}^{\text{HEG}}(\mathbb{H}; \tau, z, \nu) := \text{tr}_{\ker(\bar{H})} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} u^{\bar{J}_0 - c/6} q^H \right).$$

- [W17]

Let \mathbb{H}_0 denote the maximal \mathbb{C} -vector space such that, as a representation of $\{H, J_0, \bar{J}_0\}$, for every theory in the moduli space,

$$\mathbb{H}_0 \hookrightarrow \ker(\bar{H})$$

– the **GENERIC SPACE OF STATES**;

$$\mathcal{E}_{\text{CFT}}^{\text{HEG},0}(\mathbb{H}; \tau, z, \nu) := \text{tr}_{\mathbb{H}_0} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} u^{\bar{J}_0 - c/6} q^H \right).$$

2. Prediction for the generic space of states of K3 theories

- 3 types of irreducible $N = 4$ representations \mathcal{H}_\bullet [Eguchi/Taormina/87];
 with $\chi_\bullet(\tau, z) := \text{tr}_{\mathcal{H}_\bullet} ((-1)^{J_0} y^{J_0 - c/6} q^H)$, $\mathcal{H}_\bullet^0 := \mathcal{H}_\bullet \cap \ker(H)$:
- vacuum \mathcal{H}_0 with $\chi_0(\tau, 0) = -2$; $\mathcal{H}_0^0 \cong \mathbb{C}^2$.
 - massless matter \mathcal{H}_{mm} with $\chi_{mm}(\tau, 0) = 1$; $\mathcal{H}_{mm}^0 \cong \mathbb{C}$.
 - massive matter \mathcal{H}_h ($h \in \mathbb{R}_{>0}$), $\chi_h(\tau, z) = q^h \tilde{\chi}(\tau, z)$, $\chi_h(\tau, 0) = 0$; $\mathcal{H}_h^0 = \{0\}$.

2. Prediction for the generic space of states of K3 theories

- 3 types of irreducible $N = 4$ representations \mathcal{H}_\bullet [Eguchi/Taormina/87];
 with $\chi_\bullet(\tau, z) := \text{tr}_{\mathcal{H}_\bullet} ((-1)^{J_0} y^{J_0 - c/6} q^H)$, $\mathcal{H}_\bullet^0 := \mathcal{H}_\bullet \cap \ker(H)$:
- vacuum \mathcal{H}_0 with $\chi_0(\tau, 0) = -2$; $\mathcal{H}_0^0 \cong \mathbb{C}^2$.
 - massless matter \mathcal{H}_{mm} with $\chi_{mm}(\tau, 0) = 1$; $\mathcal{H}_{mm}^0 \cong \mathbb{C}$.
 - massive matter \mathcal{H}_h ($h \in \mathbb{R}_{>0}$), $\chi_h(\tau, z) = q^h \tilde{\chi}(\tau, z)$, $\chi_h(\tau, 0) = 0$; $\mathcal{H}_h^0 = \{0\}$.

Ansatz: $\mathbb{H} = \mathcal{H}_0 \otimes \bar{\mathcal{H}}_0 \oplus 20 \mathcal{H}_{mm} \otimes \bar{\mathcal{H}}_{mm} \oplus \left(\bigoplus_{n \in \mathbb{N}_{>0}} [f_n \mathcal{H}_n \otimes \bar{\mathcal{H}}_0 \oplus \bar{f}_n \mathcal{H}_0 \otimes \bar{\mathcal{H}}_n] \right)$
 $\oplus \left(\bigoplus_{\ell \in \mathbb{N}_{>0}} [g_\ell \mathcal{H}_\ell \otimes \bar{\mathcal{H}}_{mm} \oplus \bar{g}_\ell \mathcal{H}_{mm} \otimes \bar{\mathcal{H}}_\ell] \right) \oplus \bigoplus_{h, \bar{h} \in \mathbb{R}_{>0}} k_{h, \bar{h}} \mathcal{H}_h \otimes \bar{\mathcal{H}}_{\bar{h}}$
 where all $f_n, \bar{f}_n, g_\ell, \bar{g}_\ell, k_{h, \bar{h}}$ are non-negative integers.

2. Prediction for the generic space of states of K3 theories

3 types of irreducible $N = 4$ representations \mathcal{H}_\bullet [Eguchi/Taormina/87];
with $\chi_\bullet(\tau, z) := \text{tr}_{\mathcal{H}_\bullet}((-1)^{J_0} y^{J_0 - c/6} q^H)$, $\mathcal{H}_\bullet^0 := \mathcal{H}_\bullet \cap \ker(H)$:

- vacuum \mathcal{H}_0 with $\chi_0(\tau, 0) = -2$; $\mathcal{H}_0^0 \cong \mathbb{C}^2$.
- massless matter \mathcal{H}_{mm} with $\chi_{mm}(\tau, 0) = 1$; $\mathcal{H}_{mm}^0 \cong \mathbb{C}$.
- massive matter \mathcal{H}_h ($h \in \mathbb{R}_{>0}$), $\chi_h(\tau, z) = q^h \tilde{\chi}(\tau, z)$, $\chi_h(\tau, 0) = 0$; $\mathcal{H}_h^0 = \{0\}$.

Ansatz: $\mathbb{H} = \mathcal{H}_0 \otimes \bar{\mathcal{H}}_0 \oplus 20 \mathcal{H}_{mm} \otimes \bar{\mathcal{H}}_{mm} \oplus \left(\bigoplus_{n \in \mathbb{N}_{>0}} [f_n \mathcal{H}_n \otimes \bar{\mathcal{H}}_0 \oplus \bar{f}_n \mathcal{H}_0 \otimes \bar{\mathcal{H}}_n] \right)$
 $\oplus \left(\bigoplus_{\ell \in \mathbb{N}_{>0}} [g_\ell \mathcal{H}_\ell \otimes \bar{\mathcal{H}}_{mm} \oplus \bar{g}_\ell \mathcal{H}_{mm} \otimes \bar{\mathcal{H}}_\ell] \right) \oplus \bigoplus_{h, \bar{h} \in \mathbb{R}_{>0}} k_{h, \bar{h}} \mathcal{H}_h \otimes \bar{\mathcal{H}}_{\bar{h}}$
 where all $f_n, \bar{f}_n, g_\ell, \bar{g}_\ell, k_{h, \bar{h}}$ are non-negative integers.

- For all $\ell \in \mathbb{N}_{>0}$, $a_\ell := g_\ell - 2f_\ell$ is an invariant;
 $a_\ell \geq 0$ conjectured in [Ooguri89, W00], proved in [Eguchi/Hikami09].

2. Prediction for the generic space of states of K3 theories

3 types of irreducible $N = 4$ representations \mathcal{H}_\bullet [Eguchi/Taormina/87];
with $\chi_\bullet(\tau, z) := \text{tr}_{\mathcal{H}_\bullet} \left((-1)^{J_0} y^{J_0 - c/6} q^H \right)$, $\mathcal{H}_\bullet^0 := \mathcal{H}_\bullet \cap \ker(H)$:

- vacuum \mathcal{H}_0 with $\chi_0(\tau, 0) = -2$; $\mathcal{H}_0^0 \cong \mathbb{C}^2$.
- massless matter \mathcal{H}_{mm} with $\chi_{mm}(\tau, 0) = 1$; $\mathcal{H}_{mm}^0 \cong \mathbb{C}$.
- massive matter \mathcal{H}_h ($h \in \mathbb{R}_{>0}$), $\chi_h(\tau, z) = q^h \tilde{\chi}(\tau, z)$, $\chi_h(\tau, 0) = 0$; $\mathcal{H}_h^0 = \{0\}$.

Ansatz: $\mathbb{H} = \mathcal{H}_0 \otimes \bar{\mathcal{H}}_0 \oplus 20 \mathcal{H}_{mm} \otimes \bar{\mathcal{H}}_{mm} \oplus \left(\bigoplus_{n \in \mathbb{N}_{>0}} [f_n \mathcal{H}_n \otimes \bar{\mathcal{H}}_0 \oplus \bar{f}_n \mathcal{H}_0 \otimes \bar{\mathcal{H}}_n] \right)$
 $\oplus \left(\bigoplus_{\ell \in \mathbb{N}_{>0}} [g_\ell \mathcal{H}_\ell \otimes \bar{\mathcal{H}}_{mm} \oplus \bar{g}_\ell \mathcal{H}_{mm} \otimes \bar{\mathcal{H}}_\ell] \right) \oplus \bigoplus_{h, \bar{h} \in \mathbb{R}_{>0}} k_{h, \bar{h}} \mathcal{H}_h \otimes \bar{\mathcal{H}}_{\bar{h}}$
 where all $f_n, \bar{f}_n, g_\ell, \bar{g}_\ell, k_{h, \bar{h}}$ are non-negative integers.

- For all $\ell \in \mathbb{N}_{>0}$, $a_\ell := g_\ell - 2f_\ell$ is an invariant;
 $a_\ell \geq 0$ conjectured in [Ooguri89, W00], proved in [Eguchi/Hikami09].
- $\mathbb{H}_{\min} := \mathcal{H}_0 \otimes \bar{\mathcal{H}}_0 \oplus 20 \mathcal{H}_{mm} \otimes \bar{\mathcal{H}}_{mm} \oplus \bigoplus_{\ell \in \mathbb{N}_{>0}} a_\ell \mathcal{H}_\ell \otimes \bar{\mathcal{H}}_{mm} \subset \mathbb{H}_0$

2. Prediction for the generic space of states of K3 theories

3 types of irreducible $N = 4$ representations \mathcal{H}_\bullet [Eguchi/Taormina/87];
with $\chi_\bullet(\tau, z) := \text{tr}_{\mathcal{H}_\bullet}((-1)^{J_0} y^{J_0 - c/6} q^H)$, $\mathcal{H}_\bullet^0 := \mathcal{H}_\bullet \cap \ker(H)$:

- vacuum \mathcal{H}_0 with $\chi_0(\tau, 0) = -2$; $\mathcal{H}_0^0 \cong \mathbb{C}^2$.
- massless matter \mathcal{H}_{mm} with $\chi_{mm}(\tau, 0) = 1$; $\mathcal{H}_{mm}^0 \cong \mathbb{C}$.
- massive matter \mathcal{H}_h ($h \in \mathbb{R}_{>0}$), $\chi_h(\tau, z) = q^h \tilde{\chi}(\tau, z)$, $\chi_h(\tau, 0) = 0$; $\mathcal{H}_h^0 = \{0\}$.

Ansatz: $\mathbb{H} = \mathcal{H}_0 \otimes \overline{\mathcal{H}}_0 \oplus 20 \mathcal{H}_{mm} \otimes \overline{\mathcal{H}}_{mm} \oplus \left(\bigoplus_{n \in \mathbb{N}_{>0}} [f_n \mathcal{H}_n \otimes \overline{\mathcal{H}}_0 \oplus \overline{f}_n \mathcal{H}_0 \otimes \overline{\mathcal{H}}_n] \right)$
 $\oplus \left(\bigoplus_{\ell \in \mathbb{N}_{>0}} [g_\ell \mathcal{H}_\ell \otimes \overline{\mathcal{H}}_{mm} \oplus \overline{g}_\ell \mathcal{H}_{mm} \otimes \overline{\mathcal{H}}_\ell] \right) \oplus \bigoplus_{h, \overline{h} \in \mathbb{R}_{>0}} k_{h, \overline{h}} \mathcal{H}_h \otimes \overline{\mathcal{H}}_{\overline{h}}$
 where all $f_n, \overline{f}_n, g_\ell, \overline{g}_\ell, k_{h, \overline{h}}$ are non-negative integers.

- For all $\ell \in \mathbb{N}_{>0}$, $a_\ell := g_\ell - 2f_\ell$ is an invariant;
 $a_\ell \geq 0$ conjectured in [Ooguri89, W00], proved in [Eguchi/Hikami09].
- $\mathbb{H}_{\min} := \mathcal{H}_0 \otimes \overline{\mathcal{H}}_0 \oplus 20 \mathcal{H}_{mm} \otimes \overline{\mathcal{H}}_{mm} \oplus \bigoplus_{\ell \in \mathbb{N}_{>0}} a_\ell \mathcal{H}_\ell \otimes \overline{\mathcal{H}}_{mm} \subset \mathbb{H}_0$
- [generic chiral algebra of K3 theories] = [$N = 4$ superconformal algebra]
 $\iff \forall \ell \in \mathbb{N}, \text{generically } f_\ell = 0 \iff \mathbb{H}_{\min} = \mathbb{H}_0$.

3. Topologically half-twisted σ -models - a math version?

M : a compact Calabi-Yau D -fold, $T := T^{1,0}M$;

for any bundle $E \rightarrow M$, $\Lambda_x E := \bigoplus_{k=0}^{\infty} x^k \Lambda^k E$, $S_x E := \bigoplus_{k=0}^{\infty} x^k S^k E$

Definition [Hirzebruch88, Witten88]

$$\mathbb{E}_{q,-y} := y^{-\frac{D}{2}} \Lambda_{-y} T^* \otimes \bigotimes_{n=1}^{\infty} [\Lambda_{-y} q^n T^* \otimes \Lambda_{-y^{-1} q^n} T \otimes S_{q^n} T^* \otimes S_{q^n} T]$$

$$= y^{-\frac{D}{2}} \bigoplus_{\ell, m} q^{\ell} (-y)^m \mathcal{T}_{\ell, m}$$

$$\mathcal{E}(M; \tau, z) := y^{-\frac{D}{2}} \sum_{\ell, m} q^{\ell} (-y)^m \sum_j (-1)^j \dim H^j(M, \mathcal{T}_{\ell, m})$$

is the **COMPLEX ELLIPTIC GENUS** of M .

3. Topologically half-twisted σ -models - a math version?

M : a compact Calabi-Yau D -fold, $T := T^{1,0}M$;

for any bundle $E \rightarrow M$, $\Lambda_x E := \bigoplus_{k=0}^{\infty} x^k \Lambda^k E$, $S_x E := \bigoplus_{k=0}^{\infty} x^k S^k E$

Definition [Hirzebruch88, Witten88]

$$\mathbb{E}_{q,-y} := y^{-\frac{D}{2}} \Lambda_{-y} T^* \otimes \bigotimes_{n=1}^{\infty} [\Lambda_{-y} q^n T^* \otimes \Lambda_{-y}^{-1} q^n T \otimes S_{q^n} T^* \otimes S_{q^n} T]$$

$$= y^{-\frac{D}{2}} \bigoplus_{\ell} q^{\ell} (-y)^m \mathcal{T}_{\ell,m}$$

$$\mathcal{E}(M; \tau, z) := y^{-\frac{D}{2}} \sum_{\ell,m} q^{\ell} (-y)^m \sum_j (-1)^j \dim H^j(M, \mathcal{T}_{\ell,m})$$

is the **COMPLEX ELLIPTIC GENUS** of M .

$$\mathcal{E}(K3; \tau, z) = 8 \left(\frac{\vartheta_2(\tau, z)}{\vartheta_2(\tau, 0)} \right)^2 + 8 \left(\frac{\vartheta_3(\tau, z)}{\vartheta_3(\tau, 0)} \right)^2 + 8 \left(\frac{\vartheta_4(\tau, z)}{\vartheta_4(\tau, 0)} \right)^2.$$

3. Topologically half-twisted σ -models - a math version?

M : a compact Calabi-Yau D -fold, $T := T^{1,0}M$;

for any bundle $E \rightarrow M$, $\Lambda_x E := \bigoplus_{k=0}^{\infty} x^k \Lambda^k E$, $S_x E := \bigoplus_{k=0}^{\infty} x^k S^k E$

Definition [Hirzebruch88, Witten88]

$$\mathbb{E}_{q,-y} := y^{-\frac{D}{2}} \Lambda_{-y} T^* \otimes \bigotimes_{n=1}^{\infty} [\Lambda_{-y q^n} T^* \otimes \Lambda_{-y^{-1} q^n} T \otimes S_{q^n} T^* \otimes S_{q^n} T]$$

$$= y^{-\frac{D}{2}} \bigoplus_{\ell, m} q^{\ell} (-y)^m \mathcal{T}_{\ell, m}$$

$$\mathcal{E}(M; \tau, z) := y^{-\frac{D}{2}} \sum_{\ell, m} q^{\ell} (-y)^m \sum_j (-1)^j \dim H^j(M, \mathcal{T}_{\ell, m})$$

is the **COMPLEX ELLIPTIC GENUS** of M .

$$\mathcal{E}(K3; \tau, z) = 8 \left(\frac{\vartheta_2(\tau, z)}{\vartheta_2(\tau, 0)} \right)^2 + 8 \left(\frac{\vartheta_3(\tau, z)}{\vartheta_3(\tau, 0)} \right)^2 + 8 \left(\frac{\vartheta_4(\tau, z)}{\vartheta_4(\tau, 0)} \right)^2.$$

In general, using D copies of $bc - \beta\gamma$ systems

with modes $\beta_n^{(j)}, \gamma_n^{(j)}, b_n^{(j)}, c_n^{(j)}$, $n \in \mathbb{Z}$, $j \in \{1, \dots, D\}$:

$$S_{q^n}(\mathbb{C}^D) \cong \text{Sym}(\gamma_n^{(1)}, \dots, \gamma_n^{(D)}), \quad S_{q^n}((\mathbb{C}^D)^*) \cong \text{Sym}(\beta_n^{(1)}, \dots, \beta_n^{(D)}),$$

$$\Lambda_{-y q^n}(\mathbb{C}^D) \cong \Lambda(c_n^{(1)}, \dots, c_n^{(D)}), \quad \Lambda_{-y^{-1} q^n}((\mathbb{C}^D)^*) \cong \Lambda(b_n^{(1)}, \dots, b_n^{(D)}).$$

3. The chiral de Rham complex in a nutshell

Definition [Malikov/Schechtman/Vaintrob99]

CHIRAL DE RHAM COMPLEX Ω_M^{ch} : sheaf of super-VOAs over M ,

for any holomorphic coordinate chart $U \subset M$:

$$\Omega_M^{\text{ch}}(U) := D \text{ copies of a } bc - \beta\gamma\text{-system.}$$

Theorem [Malikov/Schechtman/Vaintrob99; Borisov/Libgober00]

$H^*(M, \Omega_M^{\text{ch}})$ (sheaf cohomology) is a topological $N = 2$ superconformal VOA.

Ω_M^{ch} is filtered with associated graded $\mathbb{E}_{q,-y}$ ($q \leftrightarrow H$, $y \leftrightarrow J_0$).

3. The chiral de Rham complex in a nutshell

Definition [Malikov/Schechtman/Vaintrob99]

CHIRAL DE RHAM COMPLEX Ω_M^{ch} : sheaf of super-VOAs over M ,
for any holomorphic coordinate chart $U \subset M$:

$$\Omega_M^{\text{ch}}(U) := D \text{ copies of a } bc - \beta\gamma\text{-system.}$$

Theorem [Malikov/Schechtman/Vaintrob99; Borisov/Libgober00]

$H^*(M, \Omega_M^{\text{ch}})$ (sheaf cohomology) is a topological $N = 2$ superconformal VOA.

Ω_M^{ch} is filtered with associated graded $\mathbb{E}_{q,-y}$ ($q \leftrightarrow H, y \leftrightarrow J_0$).

Consequence:

$$\mathcal{E}(M; \tau, z) = y^{-\frac{D}{2}} \sum_j (-1)^j \underbrace{\text{tr}_{H^j(M, \Omega_M^{\text{ch}})}((-y)^{J_0} q^H)}_{\neq \text{gr-dim}(H^j(M, \mathbb{E}_{q,-y})), \text{ in general}}$$

3. The chiral de Rham complex in a nutshell

Definition [Malikov/Schechtman/Vaintrob99]

CHIRAL DE RHAM COMPLEX Ω_M^{ch} : sheaf of super-VOAs over M ,
for any holomorphic coordinate chart $U \subset M$:

$$\Omega_M^{\text{ch}}(U) := D \text{ copies of a } bc - \beta\gamma\text{-system.}$$

Theorem [Malikov/Schechtman/Vaintrob99; Borisov/Libgober00]

$H^*(M, \Omega_M^{\text{ch}})$ (sheaf cohomology) is a topological $N = 2$ superconformal VOA.

Ω_M^{ch} is filtered with associated graded $\mathbb{E}_{q,-y}$ ($q \leftrightarrow H$, $y \leftrightarrow J_0$).

Consequence:

$$\mathcal{E}(M; \tau, z) = y^{-\frac{D}{2}} \sum_j (-1)^j \underbrace{\text{tr}_{H^j(M, \Omega_M^{\text{ch}})}((-y)^{J_0} q^H)}_{\neq \text{gr-dim}(H^j(M, \mathbb{E}_{q,-y})), \text{ in general}}$$

Result [Kapustin05]

With $\ker(\bar{H}) \subset \mathbb{H}$, the topologically half-twisted sigma model on M : the large volume limit of $\ker(\bar{H})$ is $H^*(M, \Omega_M^{\text{ch}})$.

3. The chiral de Rham complex in a nutshell

Definition [Malikov/Schechtman/Vaintrob99]

CHIRAL DE RHAM COMPLEX Ω_M^{ch} : sheaf of super-VOAs over M ,
for any holomorphic coordinate chart $U \subset M$:

$$\Omega_M^{\text{ch}}(U) := D \text{ copies of a } bc - \beta\gamma\text{-system.}$$

Theorem [Malikov/Schechtman/Vaintrob99; Borisov/Libgober00]

$H^*(M, \Omega_M^{\text{ch}})$ (sheaf cohomology) is a topological $N = 2$ superconformal VOA.

Ω_M^{ch} is filtered with associated graded $\mathbb{E}_{q,-y}$ ($q \leftrightarrow H, y \leftrightarrow J_0$).

Consequence:

$$\mathcal{E}(M; \tau, z) = y^{-\frac{D}{2}} \sum_j (-1)^j \underbrace{\text{tr}_{H^j(M, \Omega_M^{\text{ch}})}((-y)^{J_0} q^H)}_{\neq \text{gr-dim}(H^j(M, \mathbb{E}_{q,-y})), \text{ in general}}$$

Result [Kapustin05]

With $\ker(\bar{H}) \subset \mathbb{H}$, the topologically half-twisted sigma model on M : the large volume limit of $\ker(\bar{H})$ is $H^*(M, \Omega_M^{\text{ch}})$. Hence $\mathcal{E}(M; \tau, z) = \mathcal{E}_{\text{CFT}}(\mathbb{H}; \tau, z)$.

3. Hodge elliptic genera

$M, \mathbb{H}, \mathbb{E}_{q,-y} = y^{-\frac{d}{2}} \bigoplus_{\ell,m} q^\ell (-y)^m \mathcal{T}_{\ell,m}$ as before, $\nu \in \mathbb{C}, u := \exp(2\pi i \nu)$.

HODGE ELLIPTIC GENERA

- **[Kachru/Tripathy16]**

$$\mathcal{E}_{\text{CFT}}^{\text{HEG}}(\mathbb{H}; \tau, z, \nu) := \text{tr}_{\ker(\bar{H})} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} u^{\bar{J}_0 - c/6} q^H \right).$$

- **[W17]**

$$\mathcal{E}_{\text{CFT}}^{\text{HEG},0}(\mathbb{H}; \tau, z, \nu) := \text{tr}_{\mathbb{H}_0} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} u^{\bar{J}_0 - c/6} q^H \right).$$

-

$$\mathcal{E}(M; \tau, z) = y^{-\frac{d}{2}} \sum_j (-1)^j \sum_{\ell,m} q^\ell (-y)^m \dim H^j(M, \mathcal{T}_{\ell,m}).$$

$$\mathcal{E}(M; \tau, z) = y^{-\frac{d}{2}} \sum_j (-1)^j \text{tr}_{H^j(M, \Omega_M^{\text{ch}})} \left((-y)^{J_0} q^H \right).$$

3. Hodge elliptic genera

$$M, \mathbb{H}, \mathbb{E}_{q,-y} = y^{-\frac{d}{2}} \bigoplus_{\ell,m} q^\ell (-y)^m \mathcal{T}_{\ell,m} \text{ as before, } \nu \in \mathbb{C}, u := \exp(2\pi i \nu).$$

HODGE ELLIPTIC GENERA

- **[Kachru/Tripathy16]**

$$\mathcal{E}_{\text{CFT}}^{\text{HEG}}(\mathbb{H}; \tau, z, \nu) := \text{tr}_{\ker(\bar{H})} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} u^{\bar{J}_0 - c/6} q^H \right).$$

- **[W17]**

$$\mathcal{E}_{\text{CFT}}^{\text{HEG},0}(\mathbb{H}; \tau, z, \nu) := \text{tr}_{\mathbb{H}_0} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} u^{\bar{J}_0 - c/6} q^H \right).$$

- **[Kachru/Tripathy16]**

$$\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu) := (uy)^{-\frac{d}{2}} \sum_j (-u)^j \sum_{\ell,m} q^\ell (-y)^m \dim H^j(M, \mathcal{T}_{\ell,m}).$$

-

$$\mathcal{E}(M; \tau, z) = y^{-\frac{d}{2}} \sum_j (-1)^j \text{tr}_{H^j(M, \Omega_M^{\text{ch}})} \left((-y)^{J_0} q^H \right).$$

3. Hodge elliptic genera

$$M, \mathbb{H}, \mathbb{E}_{q,-y} = y^{-\frac{d}{2}} \bigoplus_{\ell,m} q^\ell (-y)^m \mathcal{T}_{\ell,m} \text{ as before, } \nu \in \mathbb{C}, u := \exp(2\pi i \nu).$$

HODGE ELLIPTIC GENERA

- **[Kachru/Tripathy16]**

$$\mathcal{E}_{\text{CFT}}^{\text{HEG}}(\mathbb{H}; \tau, z, \nu) := \text{tr}_{\ker(\bar{H})} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} u^{\bar{J}_0 - c/6} q^H \right).$$

- **[W17]**

$$\mathcal{E}_{\text{CFT}}^{\text{HEG},0}(\mathbb{H}; \tau, z, \nu) := \text{tr}_{\mathbb{H}_0} \left((-1)^{J_0 - \bar{J}_0} y^{J_0 - c/6} u^{\bar{J}_0 - c/6} q^H \right).$$

- **[Kachru/Tripathy16]**

$$\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu) := (uy)^{-\frac{d}{2}} \sum_j (-u)^j \sum_{\ell,m} q^\ell (-y)^m \dim H^j(M, \mathcal{T}_{\ell,m}).$$

- **[W17]**

$$\mathcal{E}^{\text{HEG, ch}}(M; \tau, z, \nu) := (uy)^{-\frac{d}{2}} \sum_j (-u)^j \text{tr}_{H^j(M, \Omega_M^{\text{ch}})} \left((-y)^{J_0} q^H \right).$$

4. Results

Results – if M is a K3 surface and \mathbb{H} belongs to a K3 theory:

- [Kachru/Tripathy16] (using the Bochner principle):
 $\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu)$ is independent of the complex structure.

4. Results

Results – if M is a K3 surface and \mathbb{H} belongs to a K3 theory:

- [Kachru/Tripathy16] (using the Bochner principle):
 $\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu)$ is independent of the complex structure.
- [W17] (using [W00; Creutzig/Höhn14; Song16, Song17]):
 - $\mathcal{E}^{\text{HEG, ch}}(M; \tau, z, \nu)$ is independent of the complex structure, and different from $\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu)$.

4. Results

Results – if M is a K3 surface and \mathbb{H} belongs to a K3 theory:

- [Kachru/Tripathy16] (using the Bochner principle):
 $\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu)$ is independent of the complex structure.
- [W17] (using [W00; Creutzig/Höhn14; Song16, Song17]):
 - $\mathcal{E}^{\text{HEG, ch}}(M; \tau, z, \nu)$ is independent of the complex structure, and different from $\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu)$.
 - $\mathcal{E}^{\text{HEG, ch}}(M; \tau, z, \nu) = \mathcal{E}_{\text{CFT}}^{\text{HEG, 0}}(\mathbb{H}; \tau, z, \nu)$. The generic space of states is $\mathbb{H}_0 = \mathbb{H}_{\min}$ and is given by $H^*(M, \Omega_M^{\text{ch}})$ and thus agrees with the Mathieu Moonshine module of [Eguchi/Ooguri/Tachikawa10; Gannon12].

4. Results

Results – if M is a K3 surface and \mathbb{H} belongs to a K3 theory:

- [Kachru/Tripathy16] (using the Bochner principle):
 - $\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu)$ is independent of the complex structure.
- [W17] (using [W00; Creutzig/Höhn14; Song16, Song17]):
 - $\mathcal{E}^{\text{HEG, ch}}(M; \tau, z, \nu)$ is independent of the complex structure, and different from $\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu)$.
 - $\mathcal{E}^{\text{HEG, ch}}(M; \tau, z, \nu) = \mathcal{E}_{\text{CFT}}^{\text{HEG, 0}}(\mathbb{H}; \tau, z, \nu)$. The generic space of states is $\mathbb{H}_0 = \mathbb{H}_{\min}$ and is given by $H^*(M, \Omega_M^{\text{ch}})$ and thus agrees with the Mathieu Moonshine module of [Eguchi/Ooguri/Tachikawa10; Gannon12].
 - The generic chiral algebra of K3 theories in \mathcal{M}_{K3} is the $N = 4$ superconformal algebra.

4. Results

Results – if M is a K3 surface and \mathbb{H} belongs to a K3 theory:

- [Kachru/Tripathy16] (using the Bochner principle):
 $\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu)$ is independent of the complex structure.
- [W17] (using [W00; Creutzig/Höhn14; Song16, Song17]):
 - $\mathcal{E}^{\text{HEG, ch}}(M; \tau, z, \nu)$ is independent of the complex structure, and different from $\mathcal{E}^{\text{HEG}}(M; \tau, z, \nu)$.
 - $\mathcal{E}^{\text{HEG, ch}}(M; \tau, z, \nu) = \mathcal{E}_{\text{CFT}}^{\text{HEG, 0}}(\mathbb{H}; \tau, z, \nu)$. The generic space of states is $\mathbb{H}_0 = \mathbb{H}_{\min}$ and is given by $H^*(M, \Omega_M^{\text{ch}})$ and thus agrees with the Mathieu Moonshine module of [Eguchi/Ooguri/Tachikawa10; Gannon12].
 - The generic chiral algebra of K3 theories in \mathcal{M}_{K3} is the $N = 4$ superconformal algebra.

Open:

- Is any VOA structure of \mathbb{H}_0 compatible with the M_{24} -action?
- Is M_{24} generated by symmetry surfing,
 as suggested in [Taormina/W11+...]?

THANK YOU
FOR YOUR ATTENTION!